

Tidal Phase-Shifts in Atom Interferometry: Case Study of the VLBAI Hannover

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Introduction & Motivation

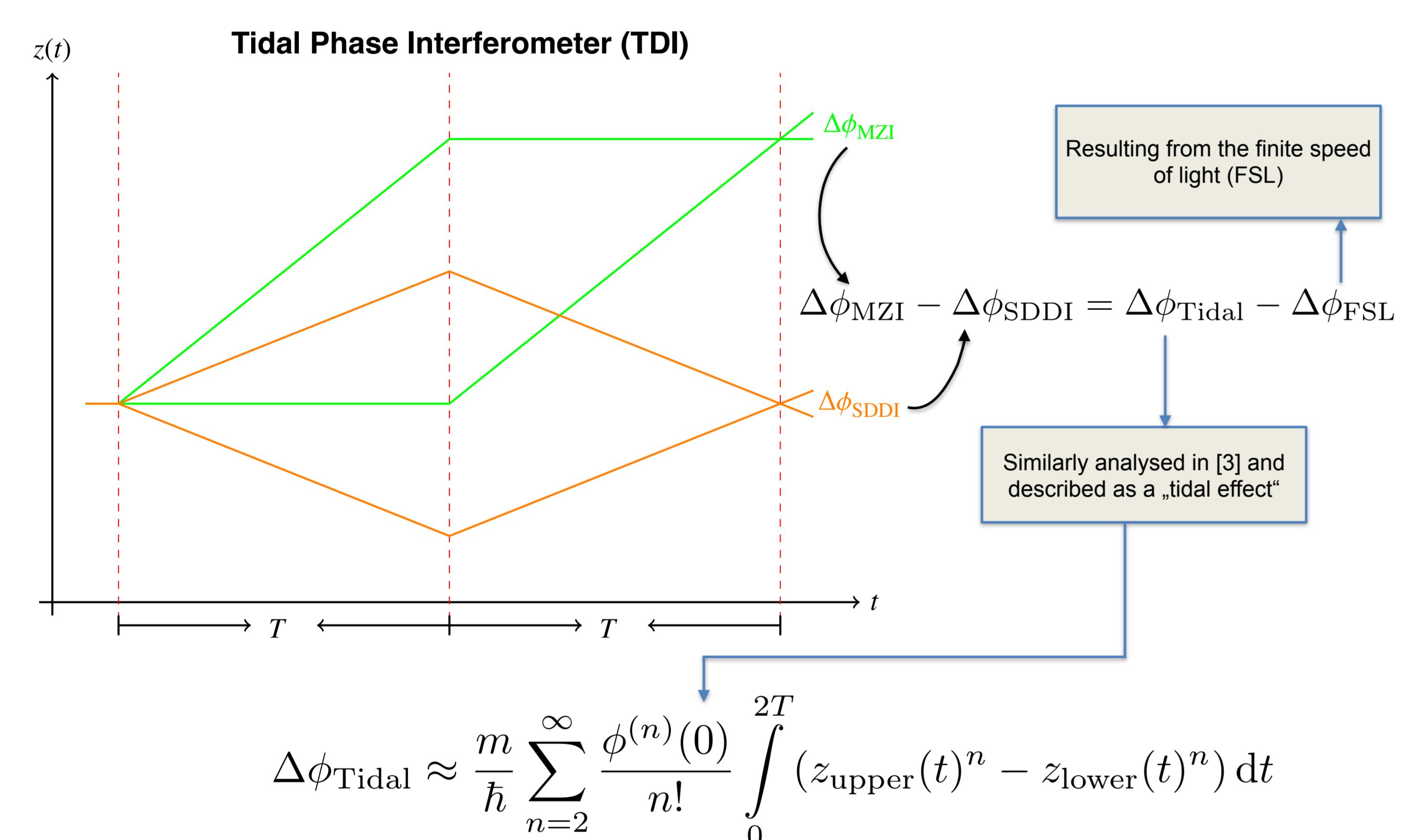
- Atom Interferometers become ever more accurate as quantum sensors, especially regarding measurements of gravity. We previously analysed how (general) relativistic effects alter interferometers (IFs) in idealised gravitational fields [1, 2].
 - Additionally, local gravitational effects have been analysed in the context of „spacetime curvature“ and gravitational tidal effects [3]. Also in recent gravitational Aharonov-Bohm-type experiments [4], macroscopic test masses have been used to generate additional sources of the gravitational field.
 - However, local gravitational effects like those can also be unintentional: Ground water variations, concrete structures, lab equipment, or even people disturb the gravitational environment. A detailed theoretical model of non-ideal gravitational fields becomes evermore important.
- ⇒ Using the gravitational field measurements of the VLBAI in Hannover, we can model future experimental setups very accurately. We do so using an open source Python algorithm.
- ⇒ We present a novel IF geometry that — dominantly — results in a phase which is connected to local gravitational field fluctuations.

Novel Interferometer Geometry for Tidal Phases

Gravitational potential expressed via its Taylor series around the origin:

$$\phi(z) = \sum_{n=0}^{\infty} \frac{\phi^{(n)}(0)}{n!} z^n$$

Consider the following differential interferometer geometry between a Mach-Zehnder Interferometer (MZI) and a Symmetric Double Diffraction Interferometer (SDDI):



This „tidal phase“ arises from the propagation phase along the quadratic term in the gravitational potential. It is therefore at least **cubic in time** and **quadratic in the photon recoil**.

The interpretation as a „tidal“ effect comes from this quadratic behaviour and is similar to the effect discussed in [3], i.e.

$$\Delta\phi_{\text{Tidal}} = -\frac{\hbar N^2 k^2 T_{zz} T^3}{2m} \quad \text{as in [3] with a grav. non-linearity } T_{zz} = \phi^{(2)}(0).$$

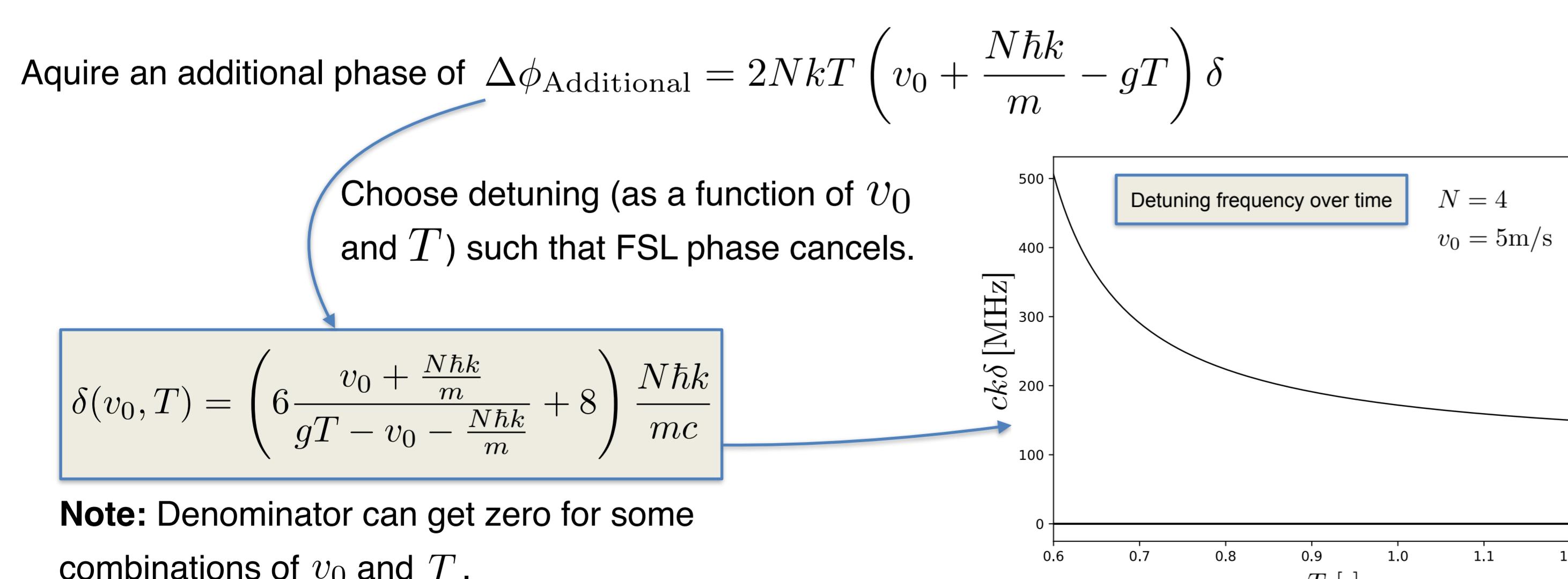
Mitigating the FSL phase shift:

The FSL phase depends on the explicit experimental setup and the type of laser interactions, i.e. being single, or two-photon interactions and the photon path lengths inside of the IF baseline [6].

Consider Bragg transitions, i.e. two light fields with individual wave vectors $k_1, -k_2$ and effective wave vector $k = k_1 + k_2$. The FSL phase is given by:

$$\Delta\phi_{\text{FSL}} = \frac{4\hbar N^2 k^2 T}{mc} \left(4gT - v_0 - \frac{N\hbar k}{m} \right)$$

Introduce a detuning of the last pulse like $\hbar k \mapsto (1 + \delta) \hbar k$ with $\delta \ll 1$.



Phase Origin - Idealized Gravitational Field

To get an understanding of how this phase originates, let us consider the idealized gravitational potential of Earth, i.e.

$$\phi(z) = gz - \frac{1}{2} \Gamma z^2 + \mathcal{O}(z^3)$$

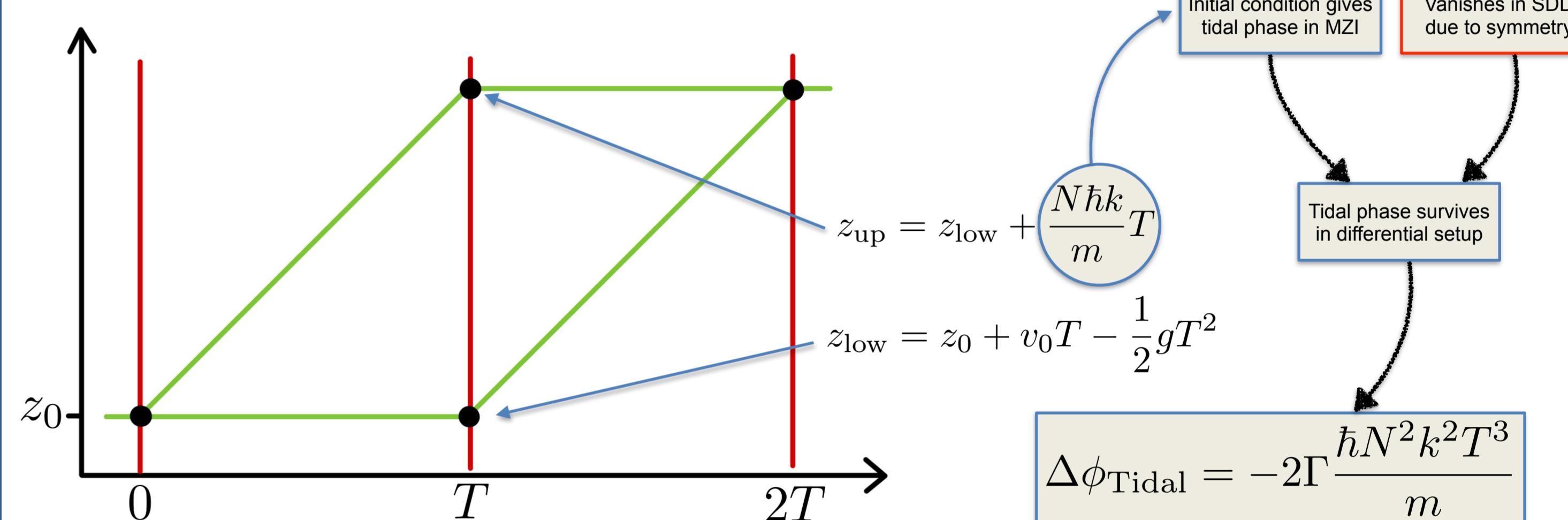
The propagation phase is given by the action functional difference along the upper and lower path:

$$\Delta\phi_{\text{Prop}} = \frac{1}{\hbar} \left(\int L(z_{\text{upper}}) dt - \int L(z_{\text{lower}}) dt \right)$$

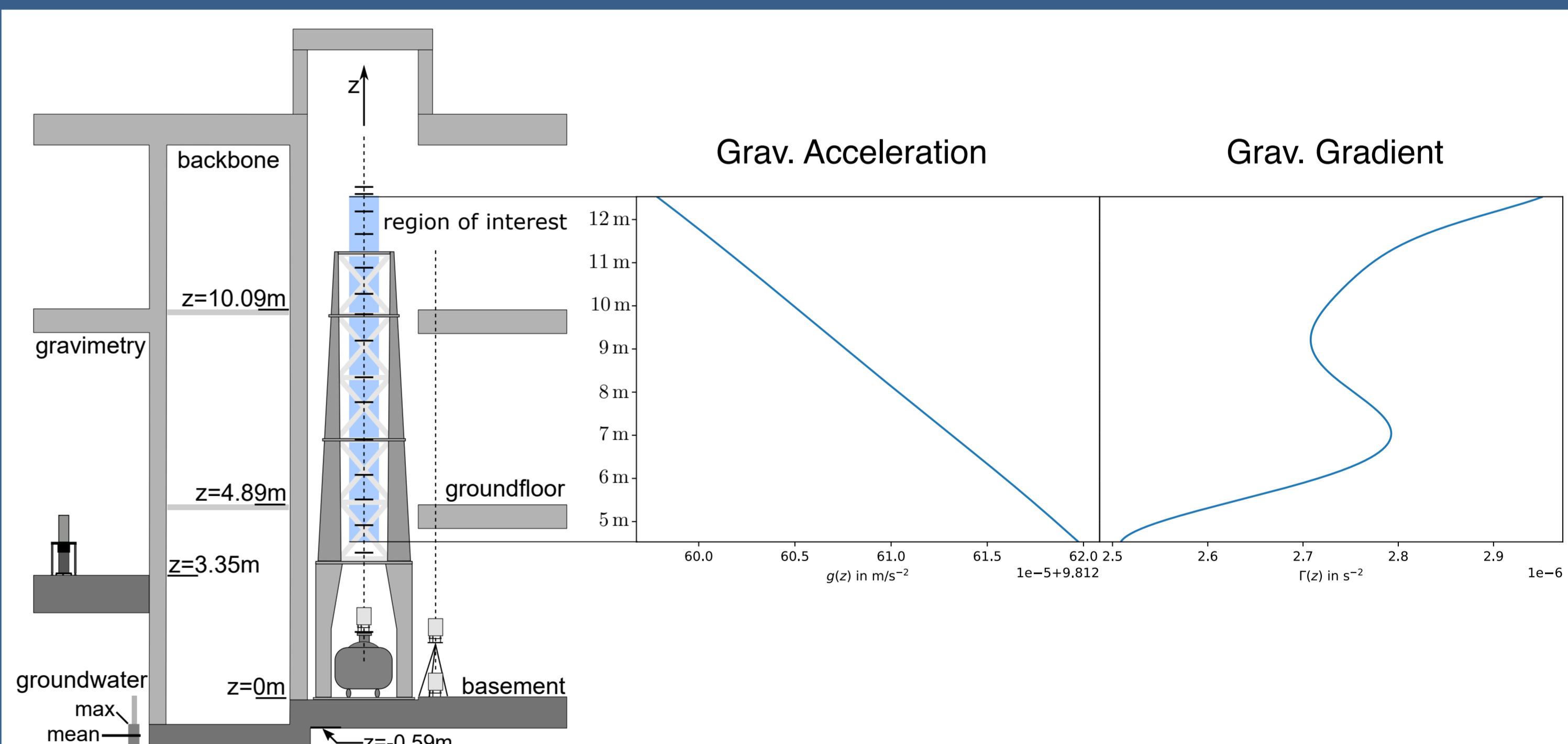
With the Lagrangian corresponding to the COM Hamiltonian

$$L(z) = \frac{m}{2} \dot{z}^2 - mgz + \frac{m}{2} \Gamma z^2$$

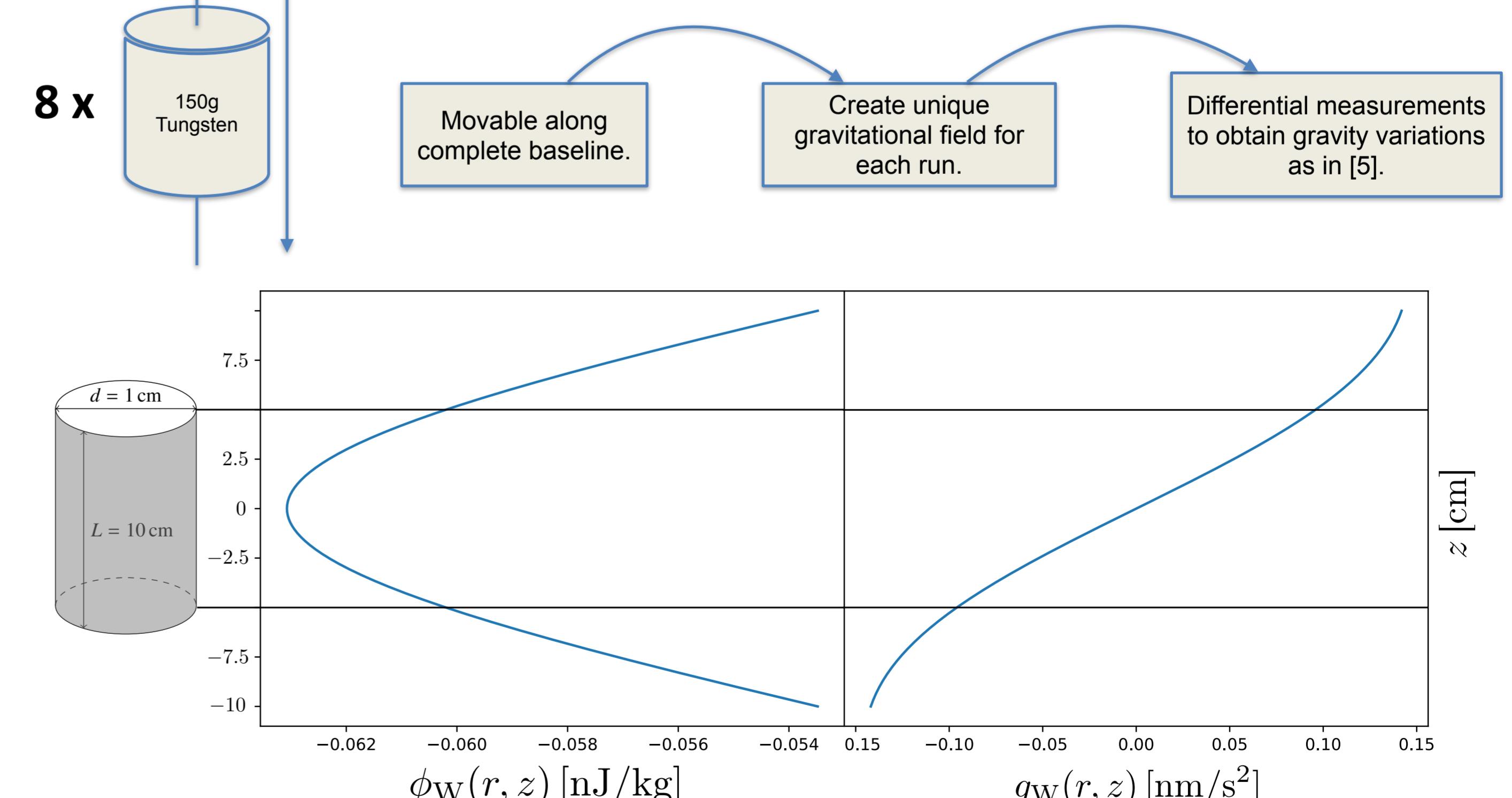
Consider the time interval $[T, 2T]$ in the propagation phase:



Real Gravitational Field: VLBAI Hannover



Additional gravitational field:



References

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- [4] C. Overstreet, P. Asenbaum, J. Curti, M. Kim, M. Kasevich, *Observation of a gravitational Aharonov-Bohm effect*, Science **375** 6577 (2022)
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- [6] Y. Tan, C.-G. Shao, Z.K. Hu, *Finite-speed-of-light perturbation in atom gravimeters*, Phys. Rev. A **94**, 013612, 2016

