

# Tidal Phase-Shifts in Atom Interferometry: Case Study of the VLBAI Hannover

Michael Werner, Dennis Schlippert, Naceur Gaaloul, Ernst Rasel and Klemens Hammerer  
Leibniz University Hannover — Germany

## Introduction & Motivation

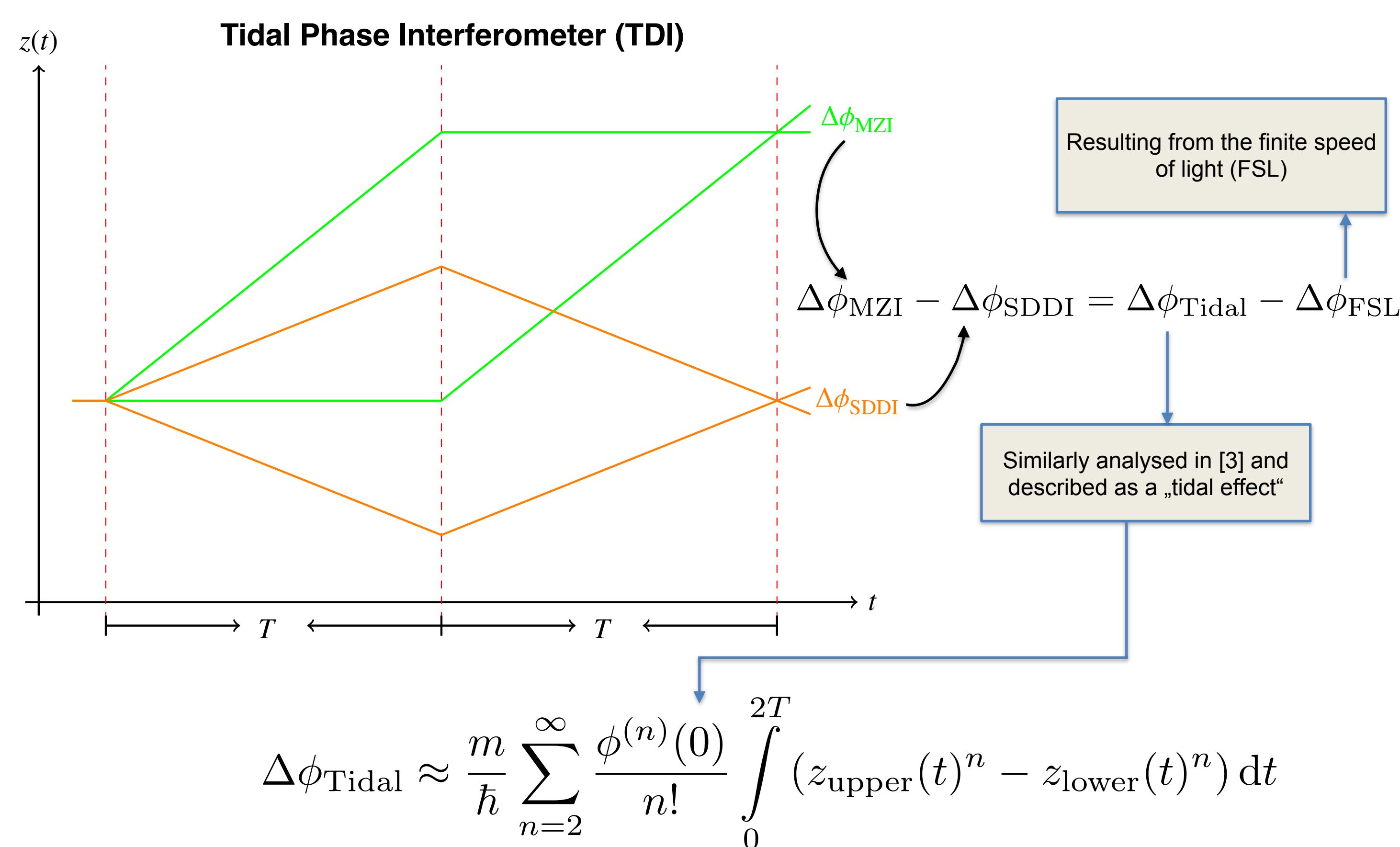
- Atom Interferometers become ever more accurate as quantum sensors, especially regarding measurements of gravity. We previously analysed how (general) relativistic effects alter interferometers (IFs) in idealised gravitational fields [1, 2].
  - Additionally, local gravitational effects have been analysed in the context of „spacetime curvature“ and gravitational tidal effects [3]. Also in recent gravitational Aharonov-Bohm-type experiments [4], macroscopic test masses have been used to generate additional sources of the gravitational field.
  - However, local gravitational effects like those can also be unintentional: Ground water variations, concrete structures, lab equipment, or even people disturb the gravitational environment. A detailed theoretical model of non-ideal gravitational fields becomes evermore important.
- ⇒ Using the gravitational field measurements of the VLBAI in Hannover, we can model future experimental setups very accurately. We do so using an open source Python algorithm.
- ⇒ We present a novel IF geometry that — dominantly — results in a phase which is connected to local gravitational field fluctuations.

## Novel Interferometer Geometry for Tidal Phases

Gravitational potential expressed via its Taylor series around the origin:

$$\phi(z) = \sum_{n=0}^{\infty} \frac{\phi^{(n)}(0)}{n!} z^n$$

Consider the following differential interferometer geometry between a Mach-Zehnder Interferometer (MZI) and a Symmetric Double Diffraction Interferometer (SDDI):



This „tidal phase“ arises from the propagation phase along the quadratic term in the gravitational potential. It is therefore at least **cubic in time** and **quadratic in the photon recoil**.

The interpretation as a „tidal“ effect comes from this quadratic behaviour and is similar to the effect discussed in [3], i.e.

$$\Delta\phi_{\text{Tidal}} = -\frac{\hbar N^2 k^2 T_{zz} T^3}{2m} \text{ as in [3] with a grav. non-linearity } T_{zz} = \phi^{(2)}(0).$$

## Mitigating the FSL phase shift:

The FSL phase depends on the explicit experimental setup and the type of laser interactions, i.e. being single, or two-photon interactions and the photon path lengths inside of the IF baseline [6].

Consider Bragg transitions, i.e. two light fields with individual wave vectors  $k_1, -k_2$  and effective wave vector  $k = k_1 + k_2$ . The FSL phase is given by:

$$\Delta\phi_{\text{FSL}} = \frac{4\hbar N^2 k^2 T}{mc} \left( 4gT - v_0 - \frac{N\hbar k}{m} \right)$$

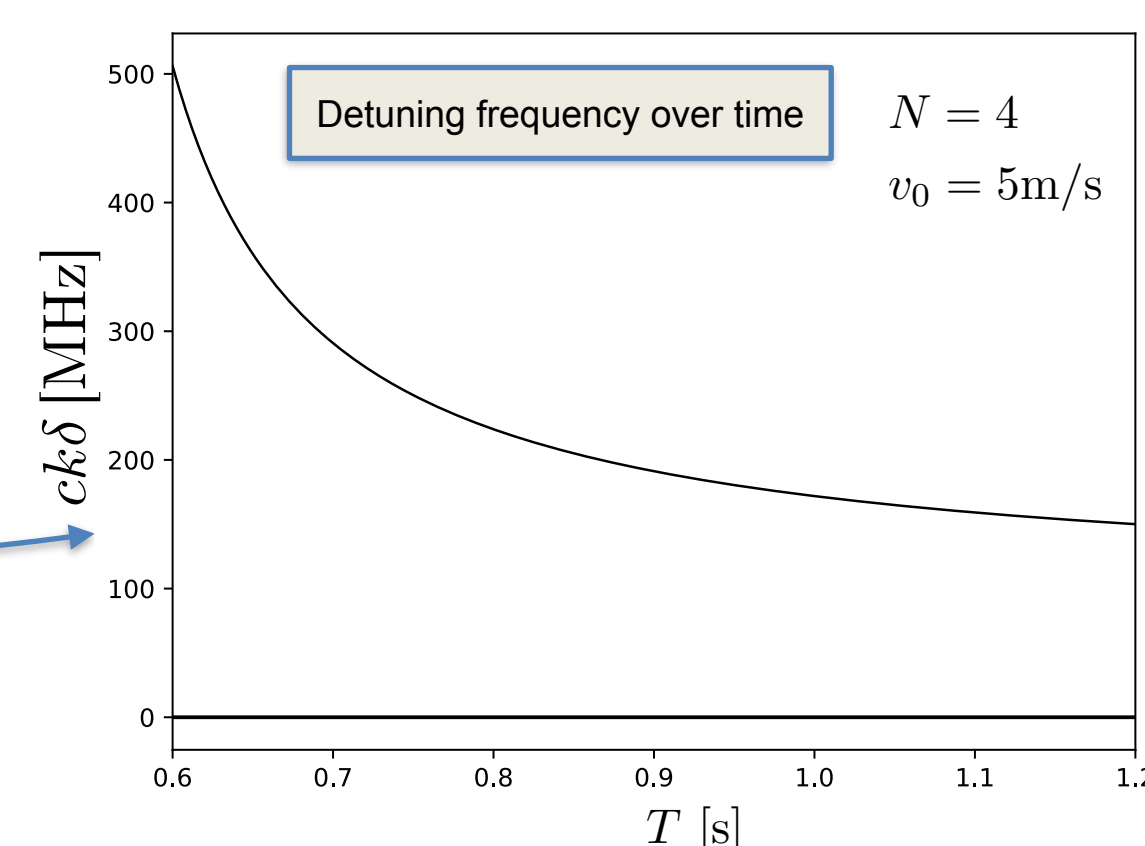
Introduce a detuning of the last pulse like  $\hbar k \mapsto (1 + \delta) \hbar k$  with  $\delta \ll 1$ .

Acquire an additional phase of  $\Delta\phi_{\text{Additional}} = 2NkT \left( v_0 + \frac{N\hbar k}{m} - gT \right) \delta$

Choose detuning (as a function of  $v_0$  and  $T$ ) such that FSL phase cancels.

$$\delta(v_0, T) = \left( 6 \frac{v_0 + \frac{N\hbar k}{m}}{gT - v_0 - \frac{N\hbar k}{m}} + 8 \right) \frac{N\hbar k}{mc}$$

**Note:** Denominator can get zero for some combinations of  $v_0$  and  $T$ .



## Phase Origin - Idealized Gravitational Field

To get an understanding of how this phase originates, let us consider the idealized gravitational potential of Earth, i.e.

$$\phi(z) = gz - \frac{1}{2}\Gamma z^2 + \mathcal{O}(z^3)$$

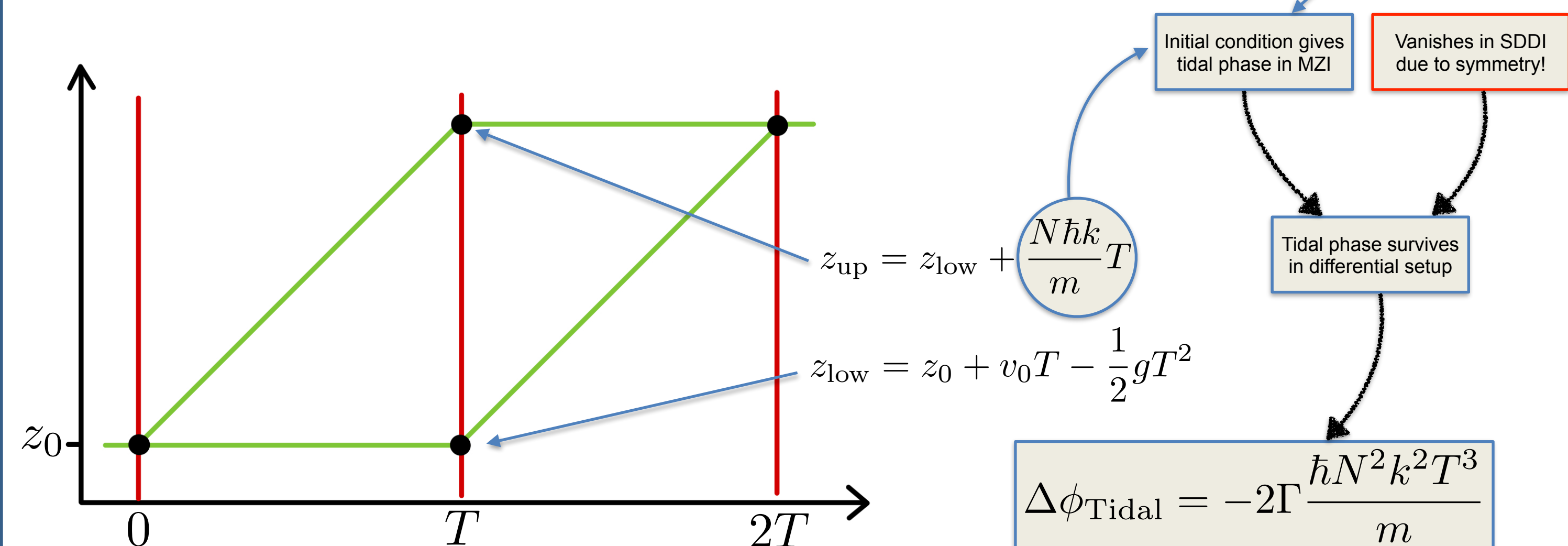
The propagation phase is given by the action functional difference along the upper and lower path:

$$\Delta\phi_{\text{Prop}} = \frac{1}{\hbar} \left( \int L(z_{\text{upper}}) dt - \int L(z_{\text{lower}}) dt \right)$$

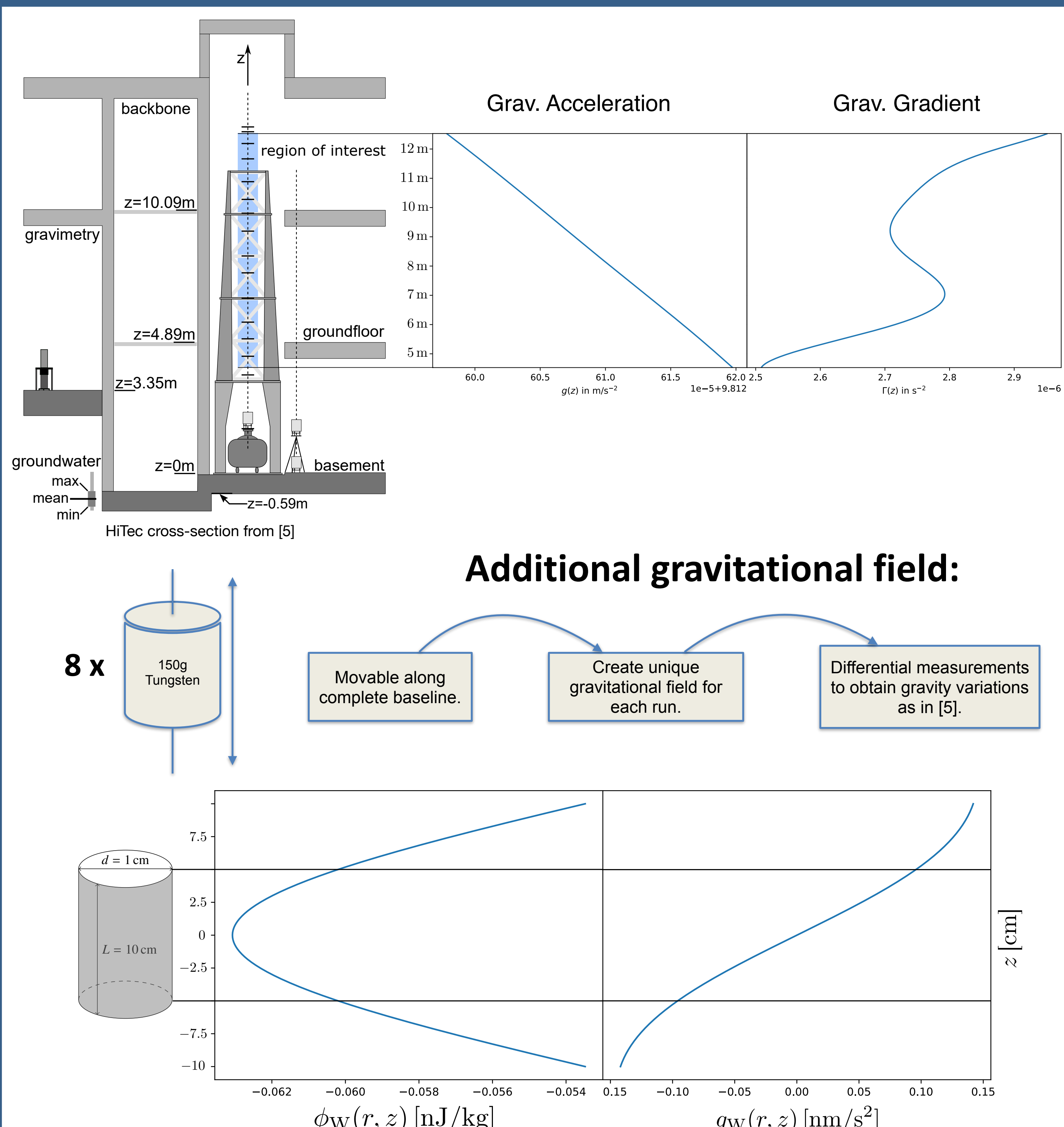
With the Lagrangian corresponding to the COM Hamiltonian

$$L(z) = \frac{m}{2} \dot{z}^2 - mgz + \frac{m}{2} \Gamma z^2$$

Consider the time interval  $[T, 2T]$  in the propagation phase:



## Real Gravitational Field: VLBAI Hannover



## References

- [1] M. Werner, P. Schwartz, J.-N. Kirsten-Siemß, N. Gaaloul, D. Giulini, K. Hammerer, *Atom interferometers in weakly curved spacetimes using Bragg diffraction and Bloch oscillations*, Phys. Rev. D **109**, 022008, 2024
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- [3] P. Asenbaum, C. Overstreet, T. Kovachy, D. Brown, J. Hogan, M. Kasevich, *Phase Shift in an Atom Interferometer due to Spacetime Curvature across its Wave Function*, Phys. Rev. Lett. **118**, 183602, 2017

- [4] C. Overstreet, P. Asenbaum, J. Curti, M. Kim, M. Kasevich, *Observation of a gravitational Aharonov-Bohm effect*, Science **375** 6577 (2022)
- [5] M. Schilling, É. Wodey, L. Timmen, D. Tell, K. Zipfel, D. Schlippert, C. Schubert, E. Rasel, J. Müller, *Gravity field modelling for the Hannover 10 m atom interferometer*, J. Geod. **94**, 122, 2020
- [6] Y. Tan, C.-G. Shao, Z.K. Hu, *Finite-speed-of-light perturbation in atom gravimeters*, Phys. Rev. A **94**, 013612, 2016

Python code [2]

